



ATOMISTIC SIMULATIONS IN PHYSICAL METALLURGY: THE ROLE OF *AB INITIO* ELECTRONIC STRUCTURE CALCULATIONS

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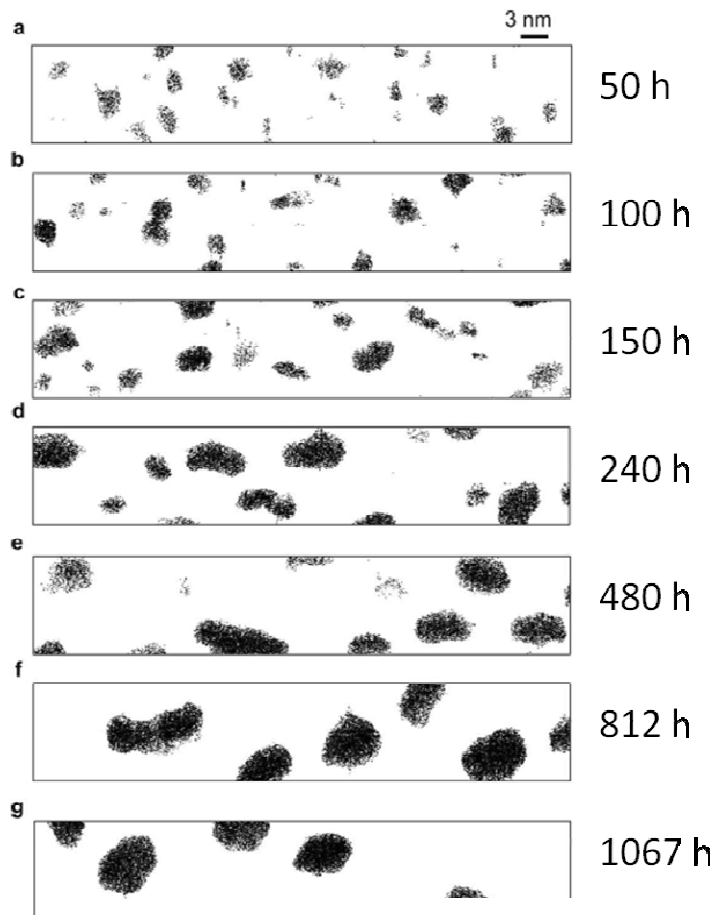
Département des Matériaux pour le Nucléaire

Service de Recherches de Métallurgie Physique

Decomposition kinetics of Fe-20%Cr during thermal aging at 500°C

Experiment

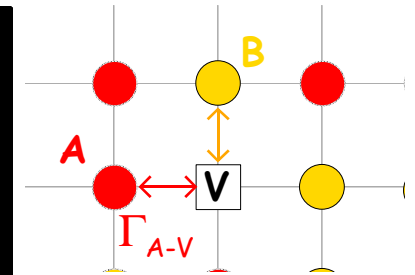
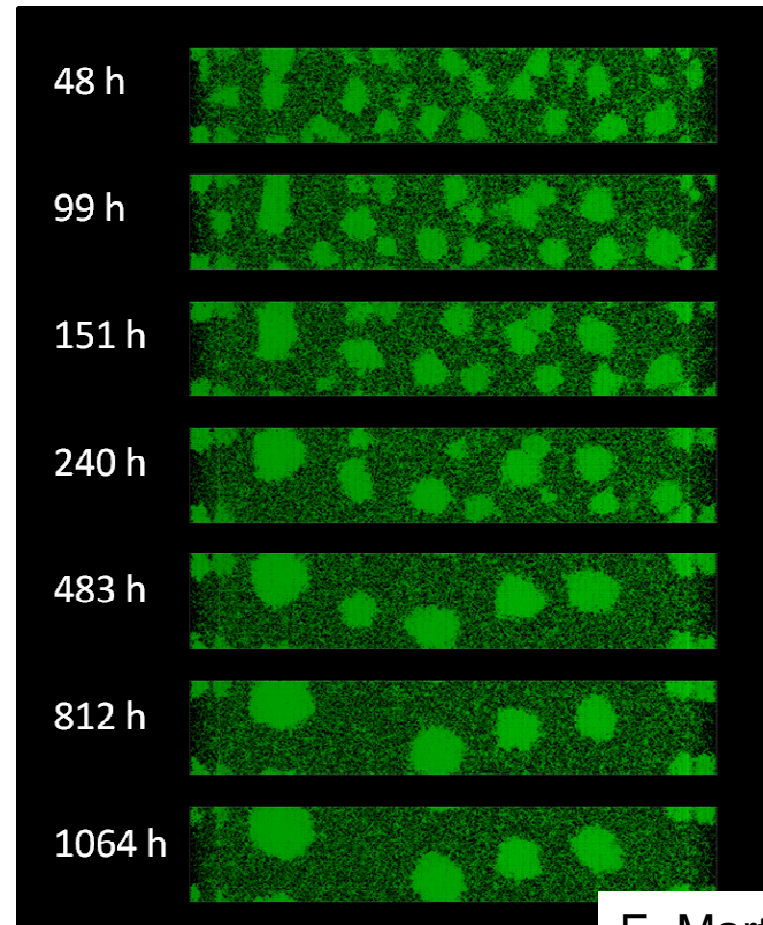
3D atom probe



Novy *et al.* (2009)

Simulation

Kinetic Monte Carlo with *ab initio* inputs



E. Martinez *et al.* (2011)

I. Introduction

II. Hierarchy of interatomic interaction models

- *Ab initio* electronic structure calculations (Density Functional Theory framework)
- Brief introduction to Tight-binding models and semi-empirical potentials

III. Different types of calculations

- Static calculations
- Structure optimization
- Energy barriers
- Quasi-harmonic calculations
- Molecular dynamics (not shown)

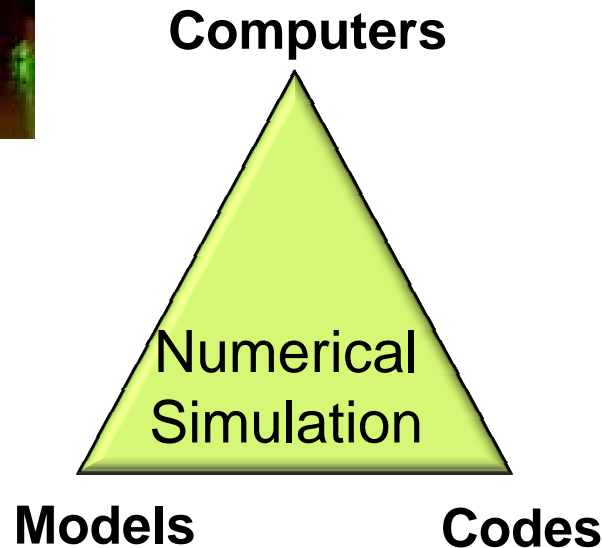
IV. A few examples

INCREASING WEIGHT OF *AB INITIO* SIMULATIONS IN MATERIALS SCIENCE

In this talk: *Ab Initio* = Density Functional Theory (DFT)



Curie, CCRT, 2 Pflops,
100 000 cores

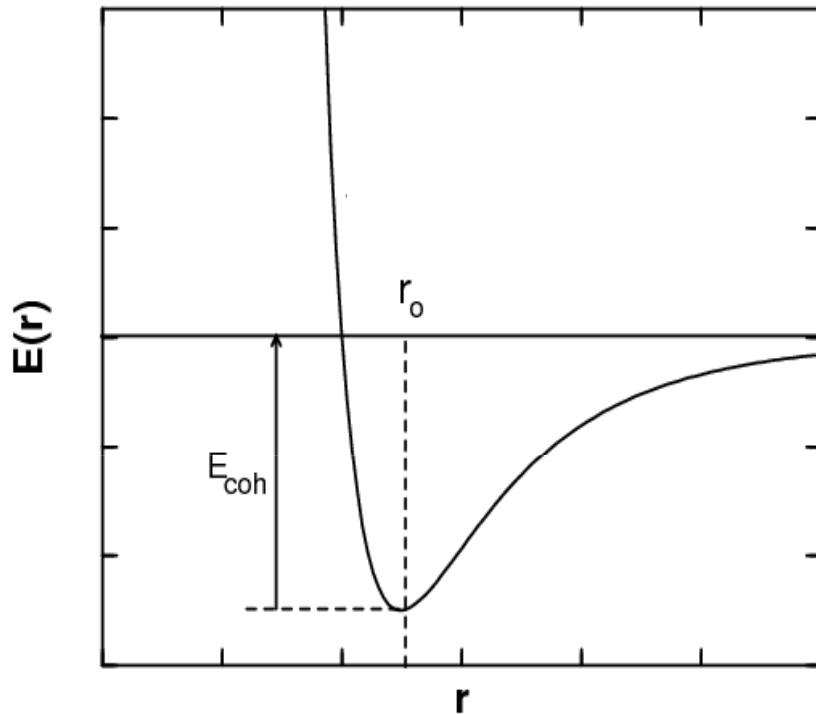


- Continuous increase of computer power (Moore's law)
- Large part of supercomputer time dedicated to materials science (in particular DFT)

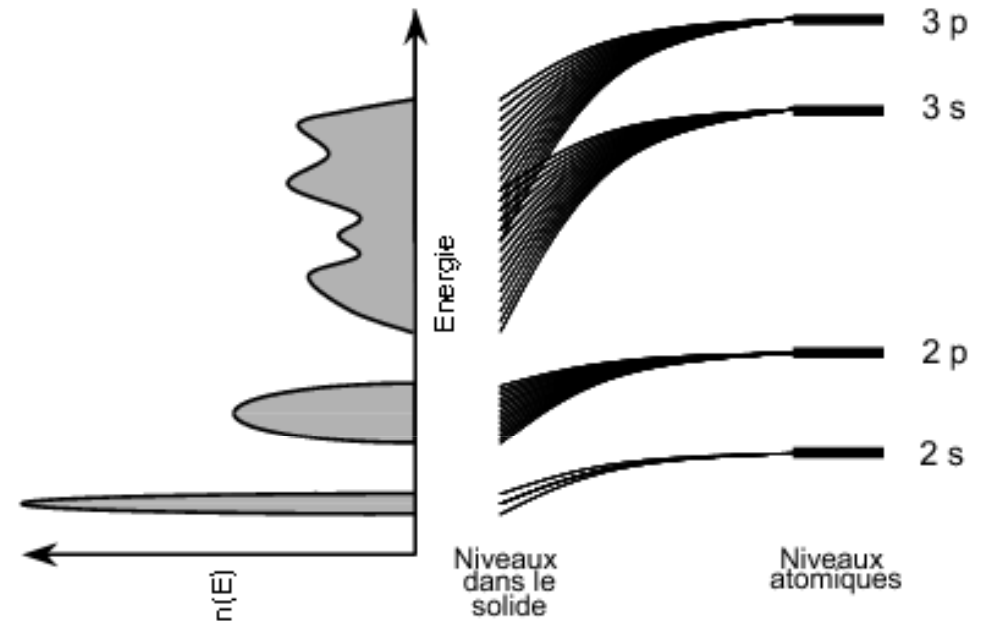
- **Efficiency of DFT** for computing quantities relevant to materials studies

- DFT codes well adapted to **massively parallel** supercomputers

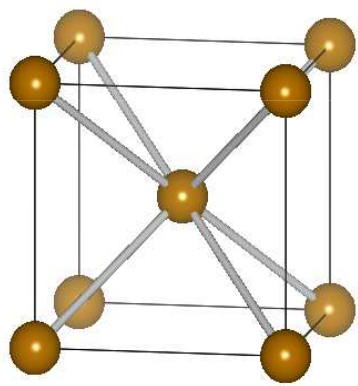
- Simulation of metals and alloys at the atomic scale
- $E(\{r_i\})$ needed
- Electronic structure origin of the binding energy
- Required accuracy in metallurgy: < 1 meV /atom



Cohesive energy: 3 to 9 eV/atom (metals)
Interatomic distance : 2 to 5 Å

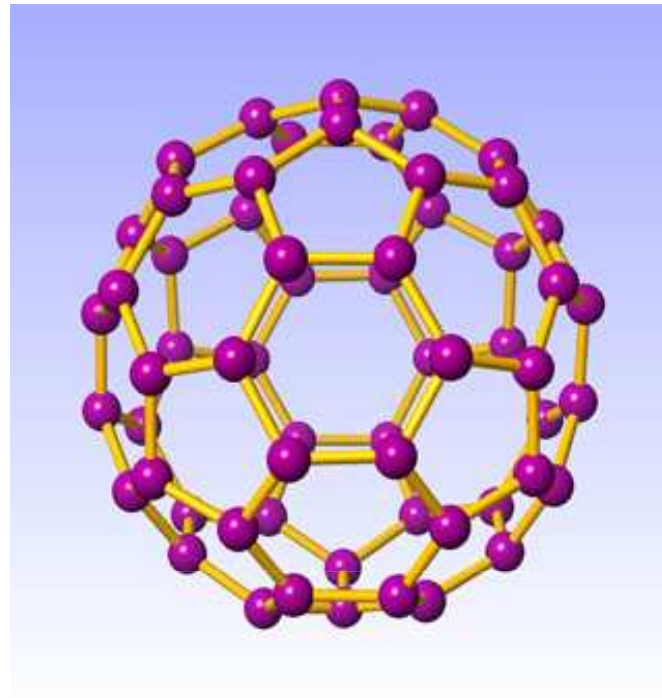


Leads to non pairwise interactions in metals
 $E_{pot} \neq \frac{1}{2} \sum_i \sum_{j \neq i} v(r_{ij})$ otherwise $E_{fv} = E_{coh}$
 instead of $\sim E_{coh}/3$

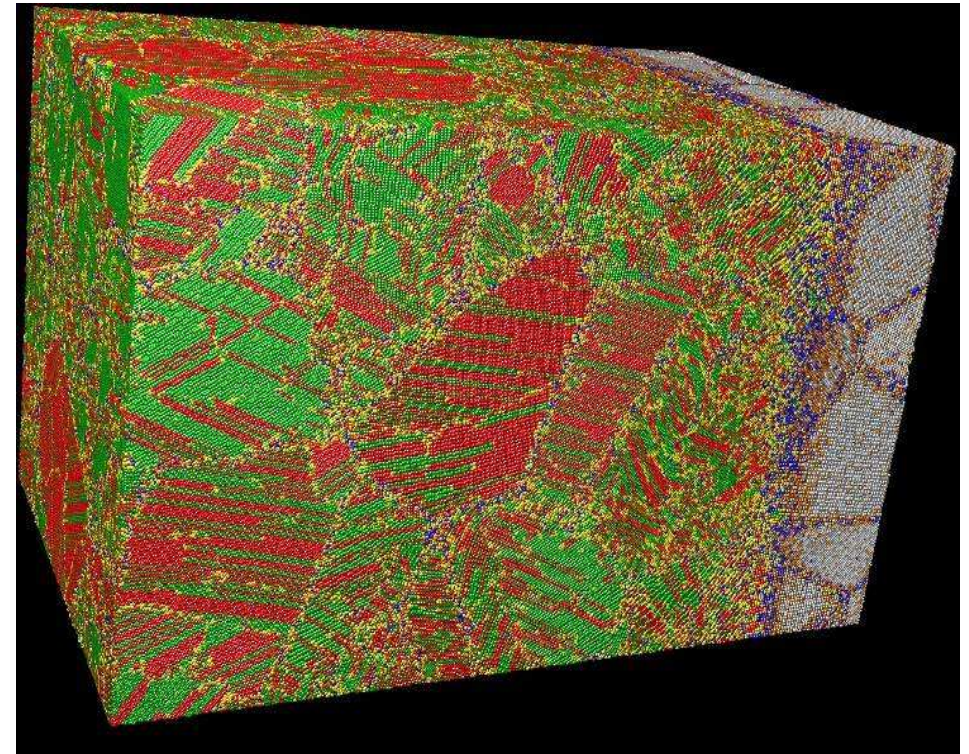


0.3 nm

1 atom (per unit cell)



1 nm
60 atoms



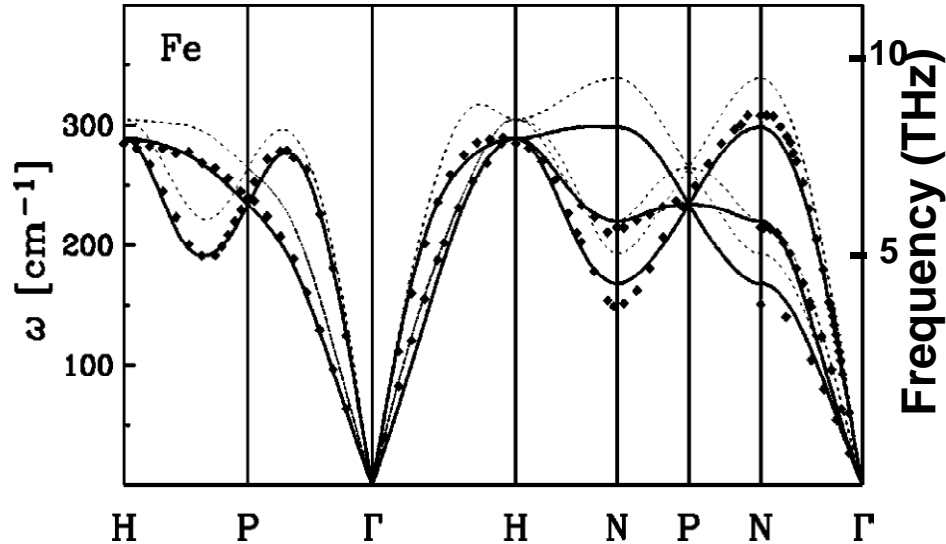
100 nm

- Example : $30 \cdot 10^6$ atoms (K. Kadau, LANL, 2007)

TIME SCALES IN ATOMISTIC SIMULATIONS

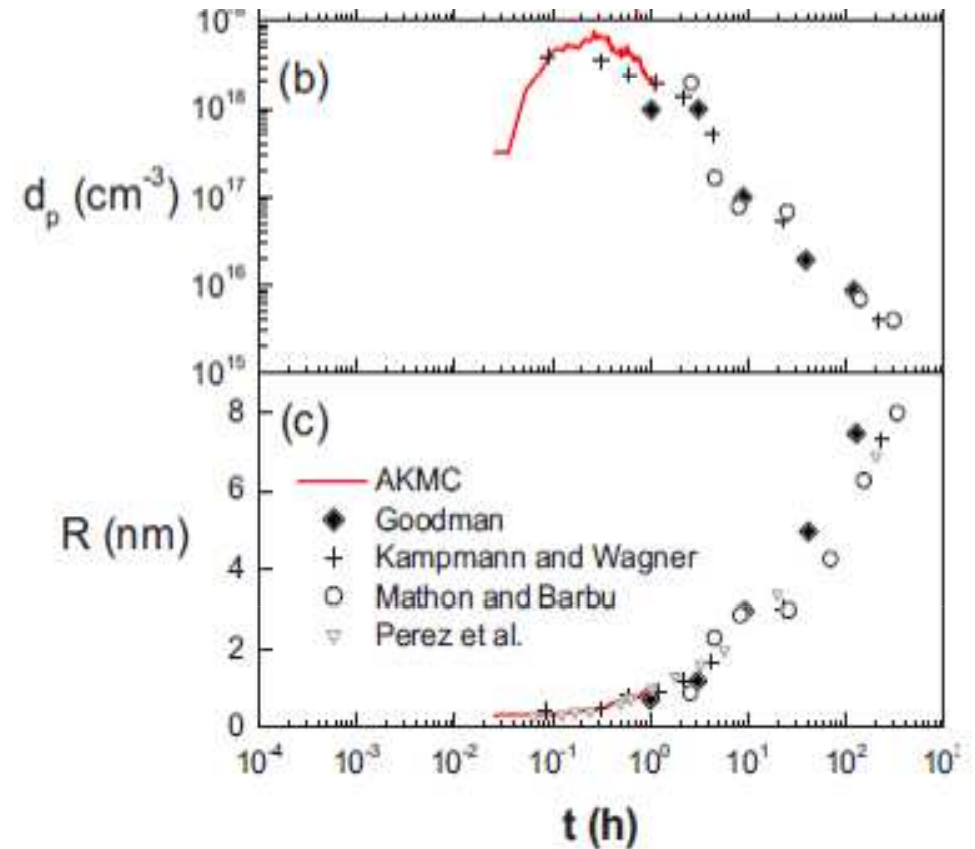
Lattice vibrations

Picosecond (10^{-12} sec)



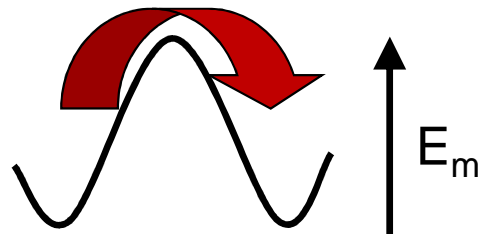
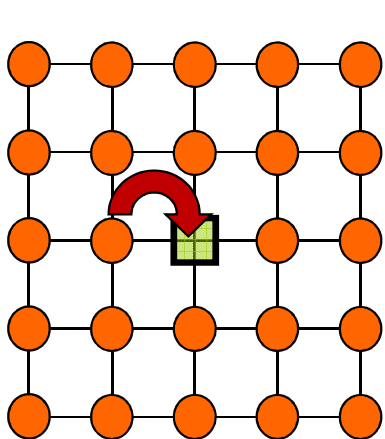
Long range diffusion

Second / hour



Atomic jump

Nanosecond (10^{-9} sec)



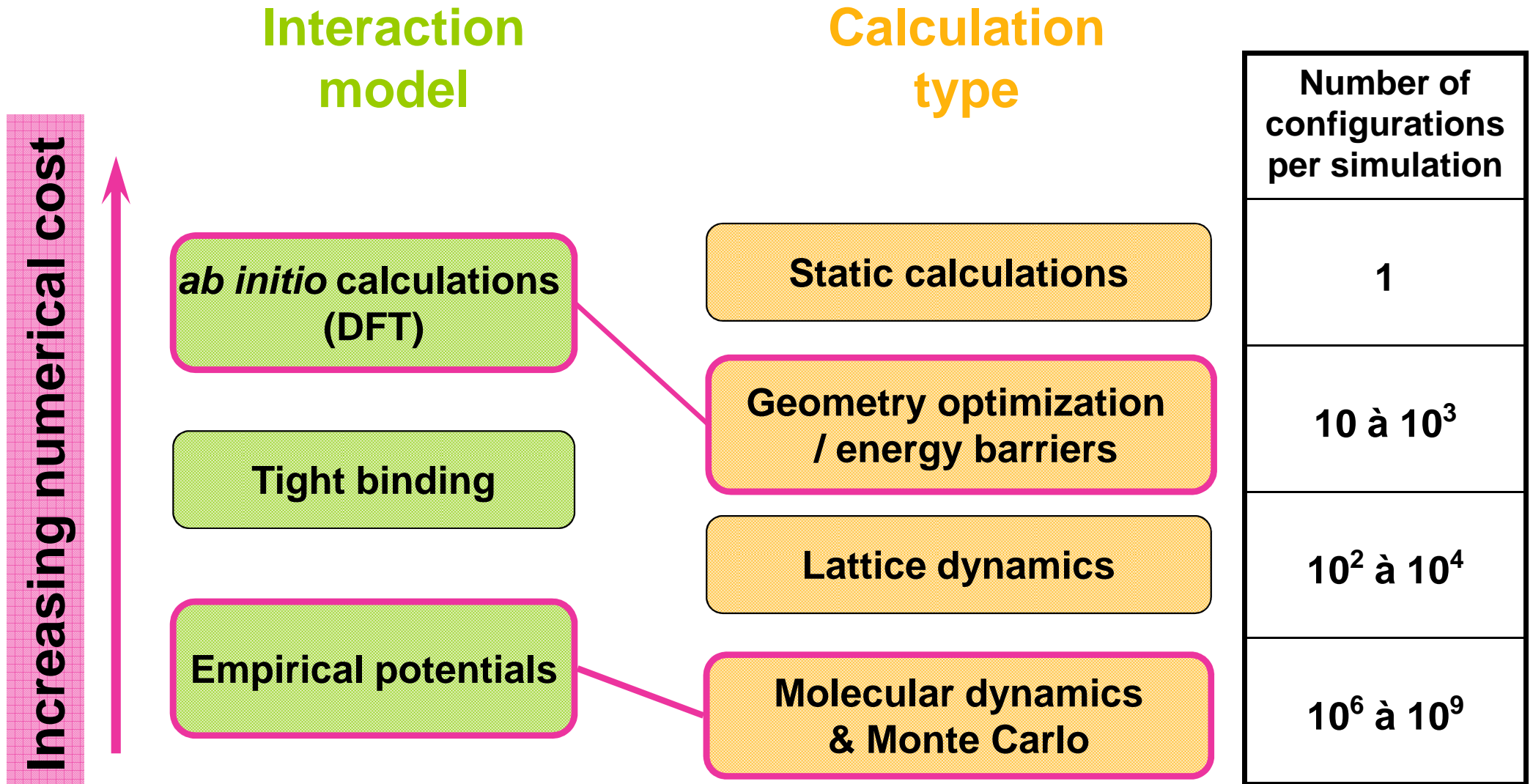
$$v = v_0 \exp(-E_m / k_B T)$$

$$v_0 = 5 \text{ THz} ; E_m = 0.5 \text{ eV}$$

$$T = 300 \text{ K}$$

Cu precipitation in Fe
Lattice kinetic Monte Carlo

(F. Soisson, CEA, 2007)



AB INITIO ELECTRONIC STRUCTURE CALCULATIONS (DFT)

- Quantum solution of a single particle approximation of Schrödinger's equation
- No adjustable parameter

- Density Functional Theory (DFT)

$$E = \min E[\rho] \quad \text{Hohenberg \& Kohn (1964)}$$

- Kohn-Sham Theory (1965)

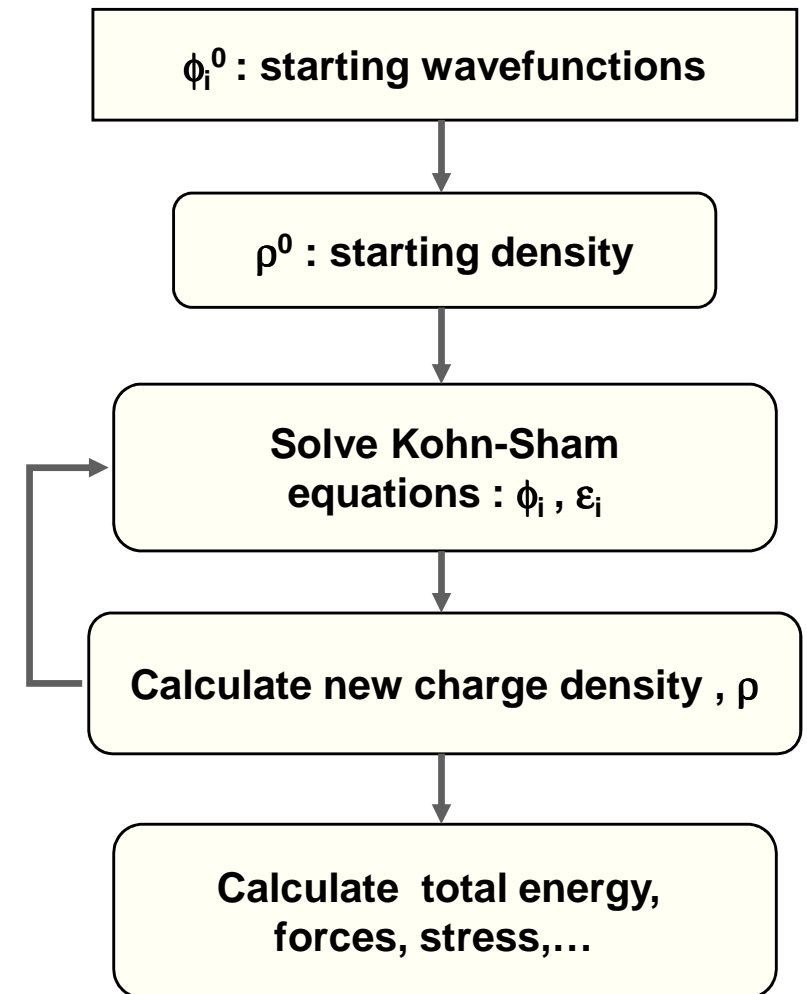
$$E[\rho] = E_{I-I} + E_{e-I} + E_H + T + E_{xc}$$

- Parametrize the particle density in terms of a set of one-electron orbitals representing a non-interacting reference system

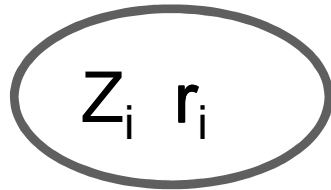
$$\left(-\frac{1}{2} \nabla^2 + v(r) + \int \frac{\rho(r')}{|r-r'|} dr' + \mu_{xc}(\rho(r)) \right) \varphi_i = \varepsilon_i \varphi_i$$

- Local Density Approximation (LDA)

- exchange-correlation energy of a homogeneous electron gas with the local density $\rho(r)$



“no adjustable parameter”



Most properties of most materials can be predicted with an accuracy of a few %:
 a , B , C_{ij} , $\omega(k)$, ΔE (structures, surfaces),
 E_{defects}

Approximation: exchange correlation functional (LDA, GGA, ...), difference gives error bar

Compromises (speed / accuracy)

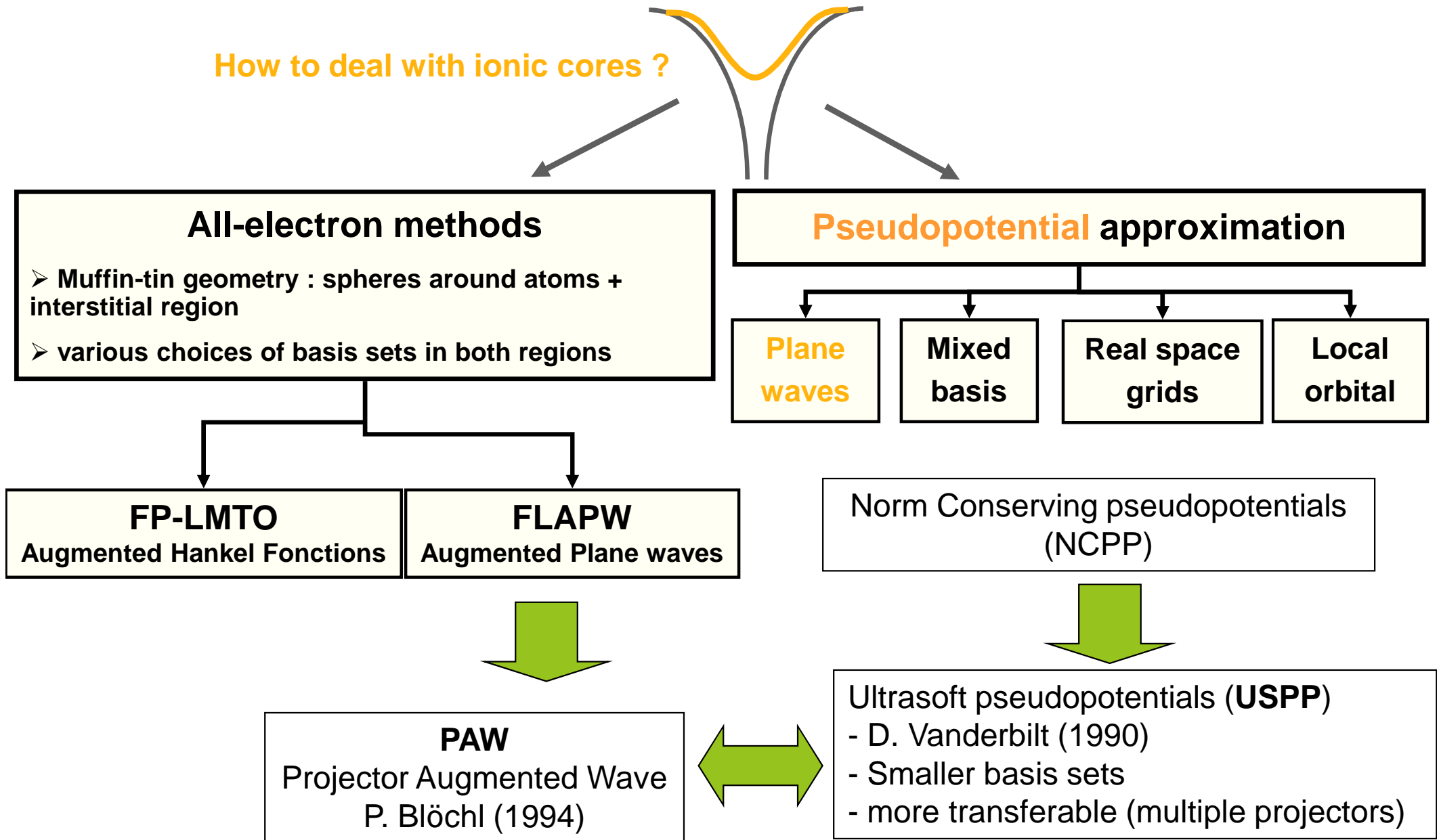
- Valence electrons ; pseudopotential
- Basis set size
- Spin polarization
- Brillouin zone integration
- Convergence criteria (electronic structure, atomic relaxation)

Limitations

- Exceptions: gap, strongly correlated electrons
- Cell size (max 100 to 1000 atoms)
- Finite temperature effects

Standard = pseudopotentials or PAW and plane waves (implemented in VASP, QuantumEspresso, ABINIT)

How to deal with ionic cores ?



DFT CODES (NON EXHAUSTIVE LIST)

Plane wave and related (real space, wavelet, etc.) methods

- **VASP**
- **CASTEP and CETEP**
- **CPMD**
- **ABINIT** - GPL
- **BigDFT** - wavelets
- **Quantum-Espresso (formerly PWscf)** - GPL
- **PEtot** - GPL
- **DACAPO** - GPL
- **Socorro** - GPL
- **DFT++** - GPL
- **Octopus** - GPL
- **Paratec**
- **DoD Planewave**
- **ACRES**
- **PARSEC** - GPL - real space, pseudopotential
- **CP2K** - GPL (mixed basis DFT)
- **GPAW** - GPL - real-space multigrid PAW

Local orbital basis codes

- **QUEST: SeqQuest** - gaussian basis pseudopotential
- **SIESTA** - numerical atom-centered basis pseudopotential
- **CRYSTAL - CSE** - gaussian basis all-electron
- **AIMPRO**
- **FPLO**
- **OpenMX** - GPL - numerical atom-centered basis PP code (Ozaki group)

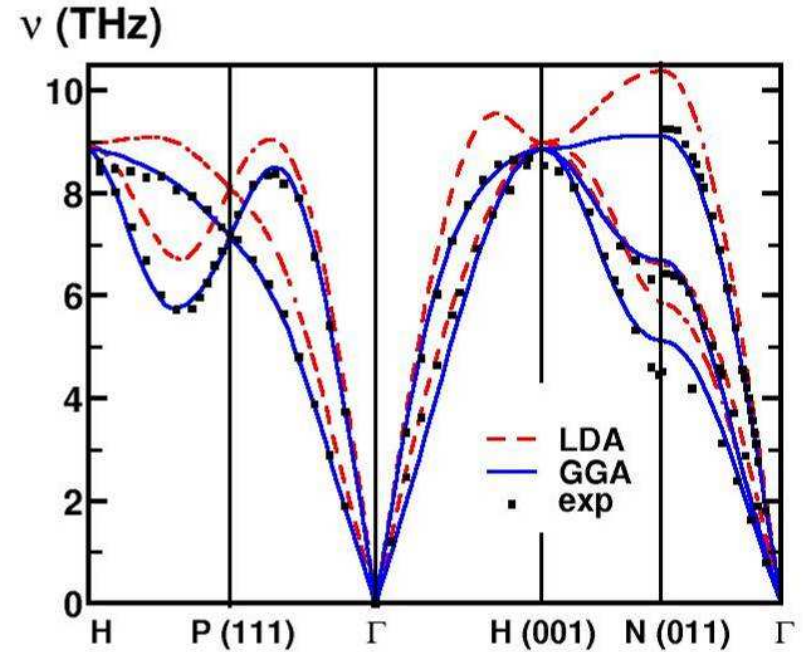
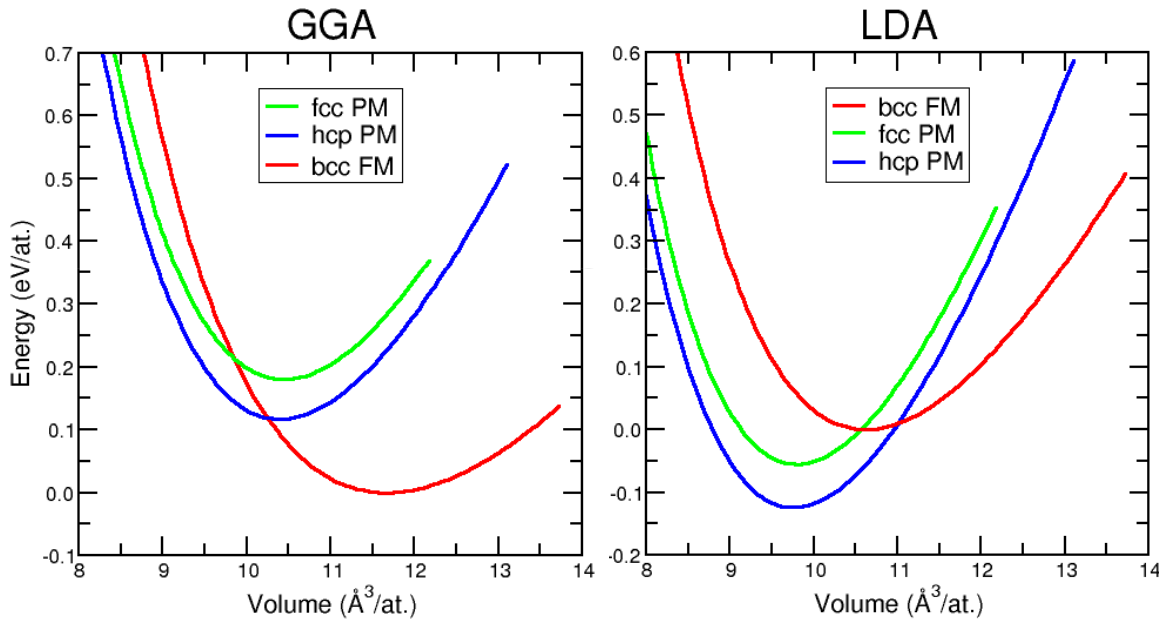
All-electron (augmented methods) codes

- **EXCITING** - GPL code for FP-LAPW
- **WIEN 2k**
- **FLAPW**
- **RSPt** - FP-LMTO

source :

http://dft.sandia.gov/Quest/DFT_codes.html

LDA vs GGA in iron



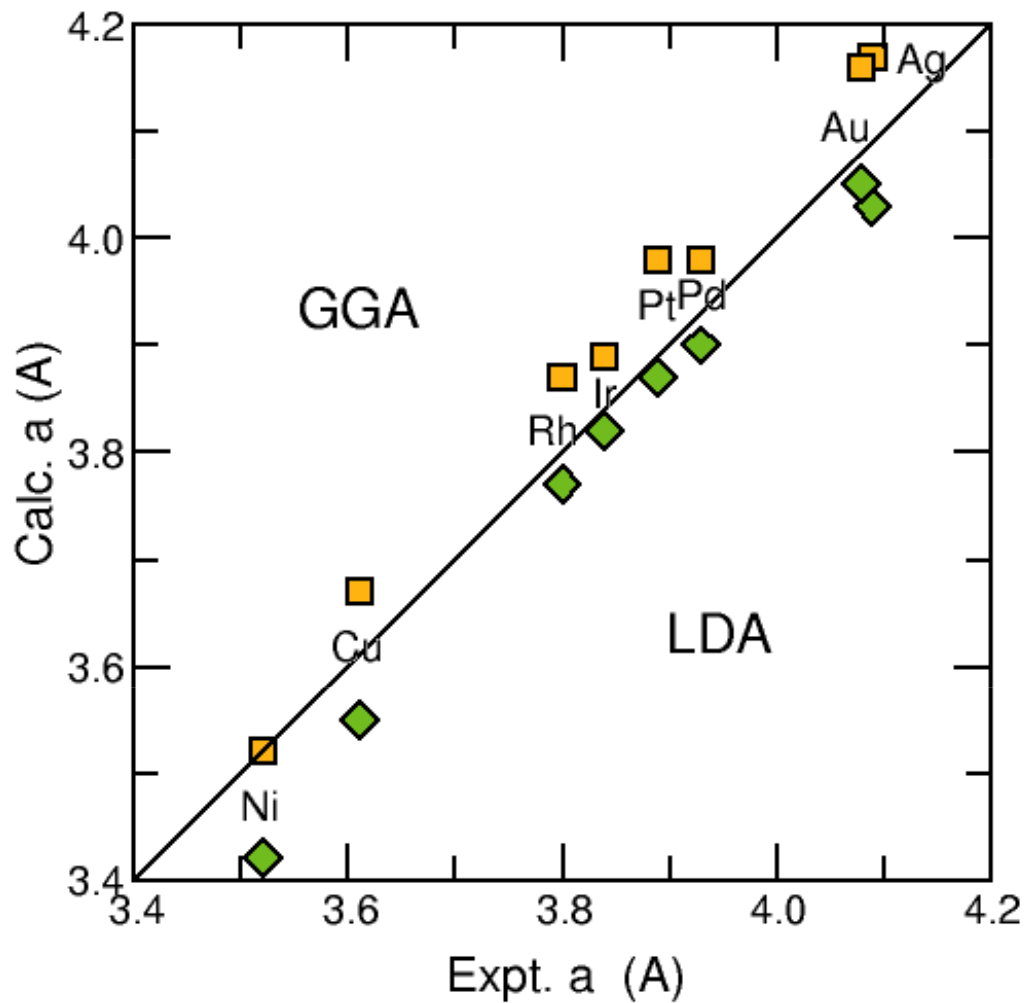
	Bulk			Vacancy	
	a (Å)	B (Mbar)	μ (μ _B)	Formation (eV)	Migration (eV)
GGA (various pseudos)	2.86	1.67 (1.52 - 1.67)	2.27 (2.26-2.38)	2.09 (1.93 - 2.09)	0.67 (0.59 - 0.67)
LDA	2.77	2.41	2.10	2.40	0.68
Exp.	2.87	1.68	2.22	2.35 ± 0.2	0.55 ± 0.05

GGA better

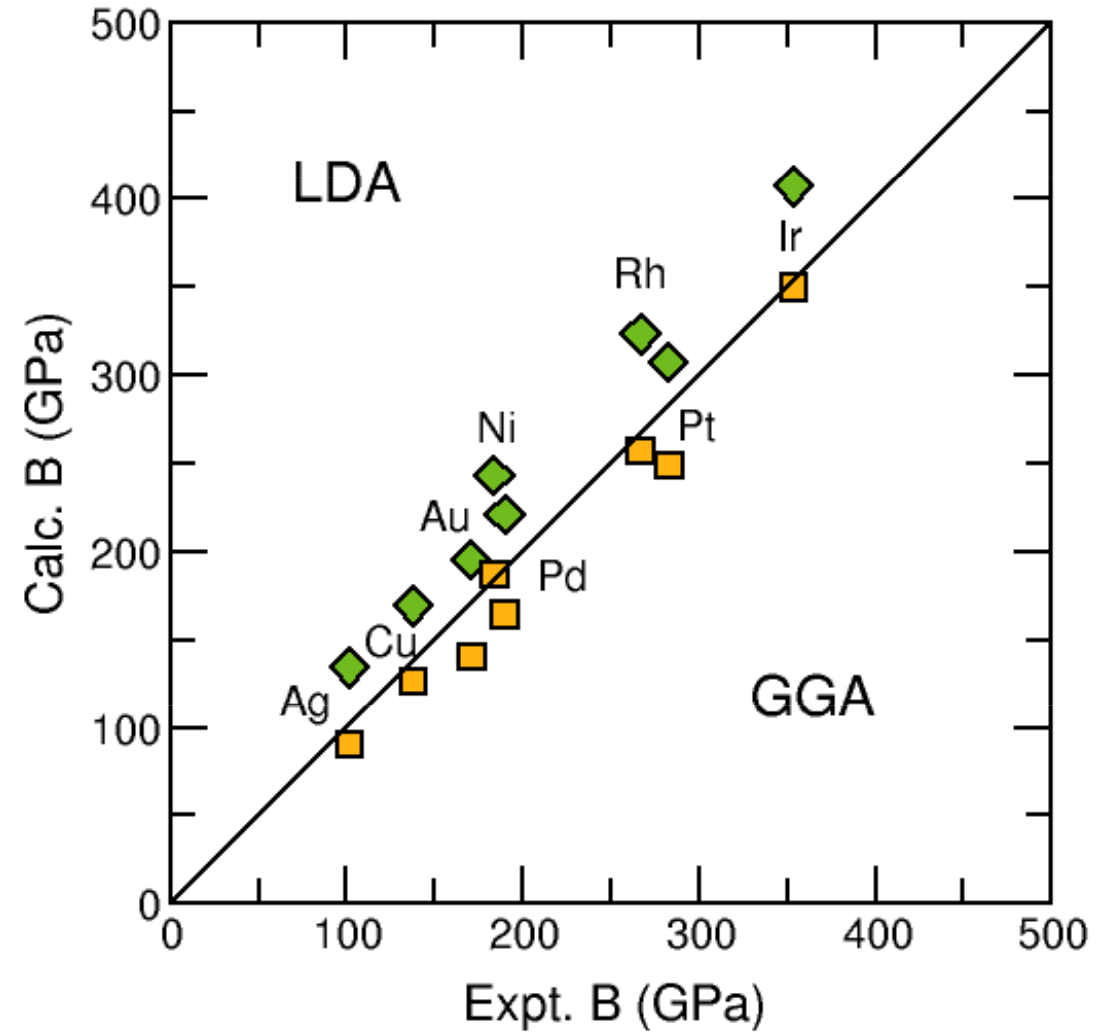
LDA better

LDA vs GGA: bulk properties of FCC metals

Lattice parameter

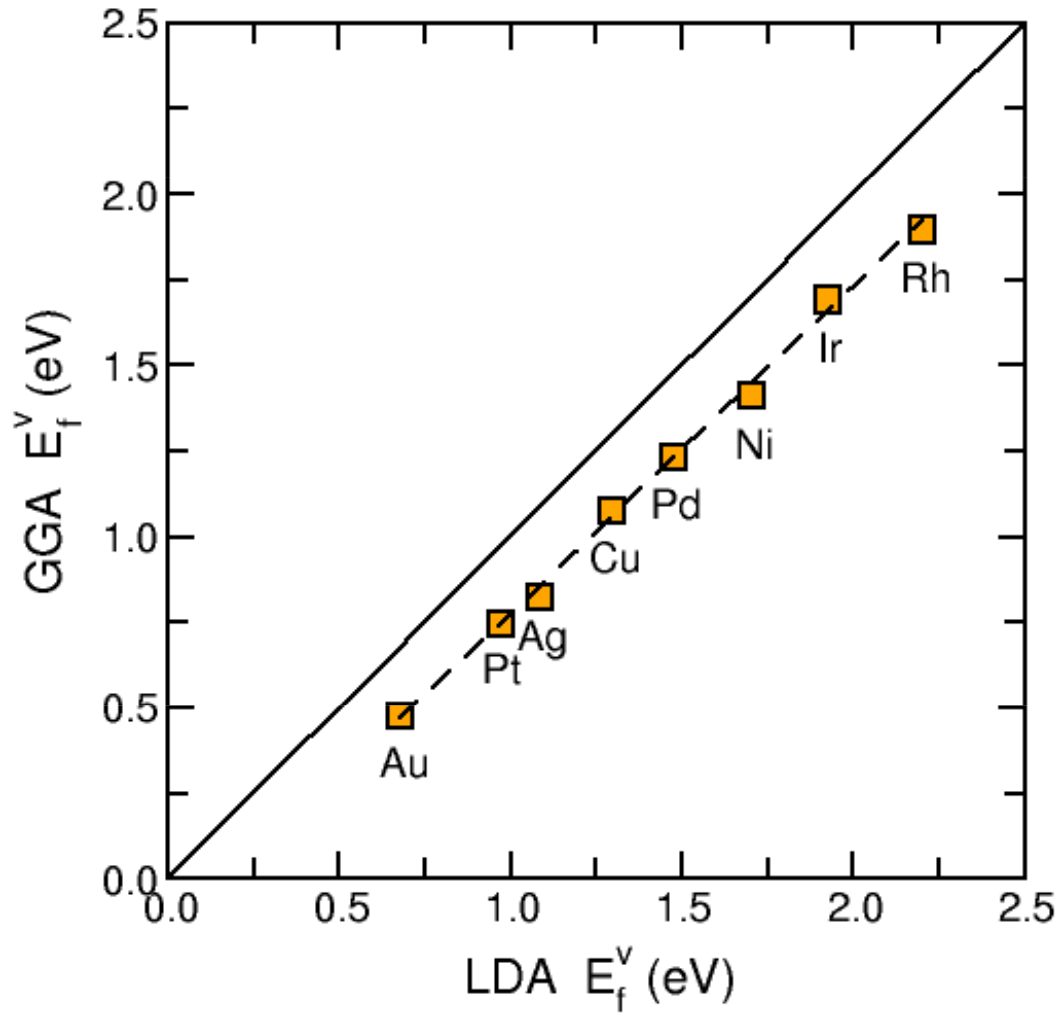


Bulk modulus



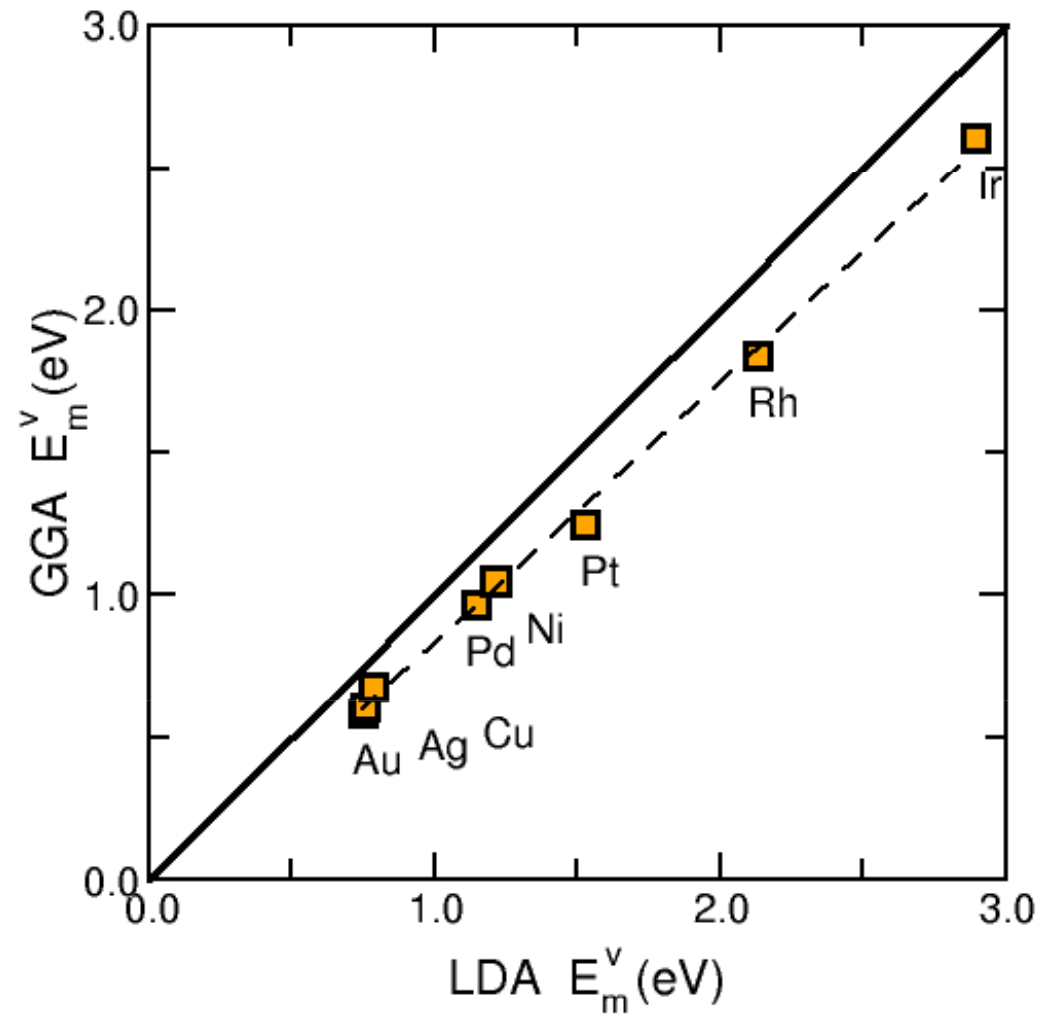
LDA underestimates a and overestimates B
GGA improves in particular for 3d metals

Vacancy formation energy



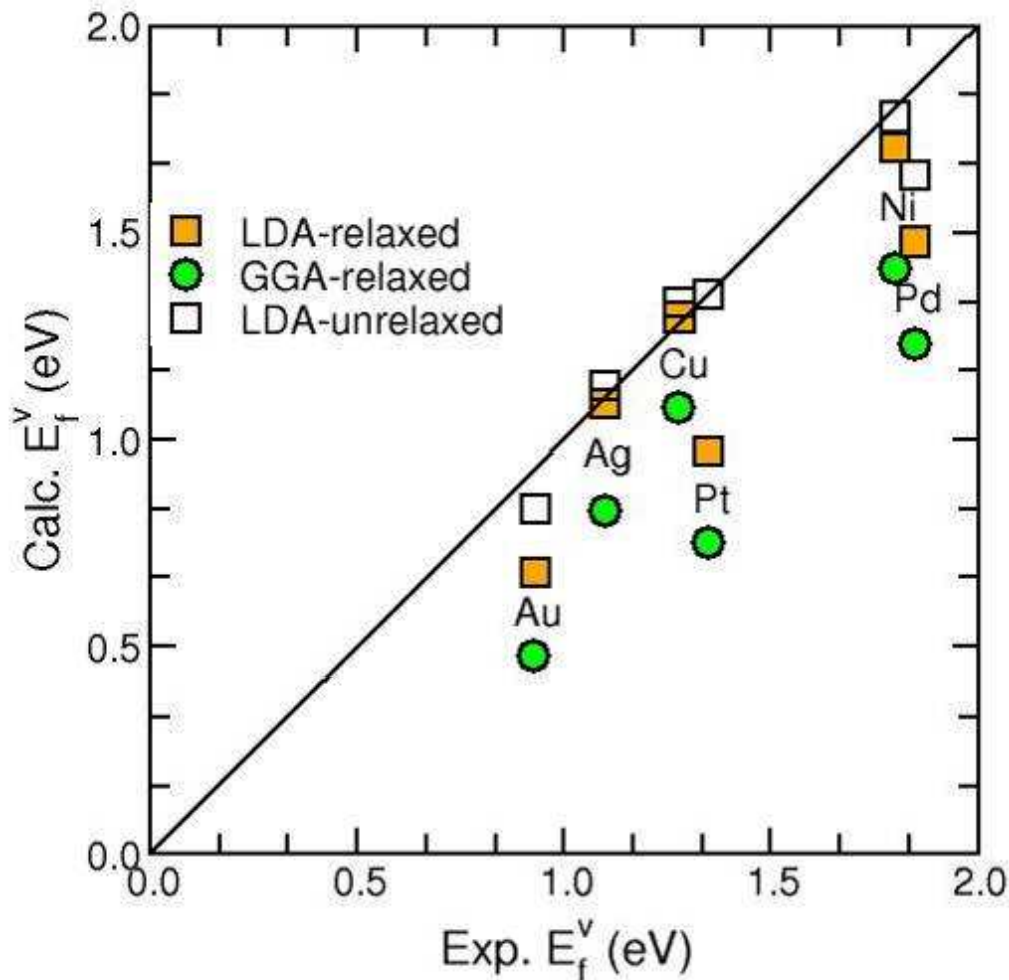
GGA values lower than LDA
by 0.25 ± 0.05 eV

Vacancy migration energy

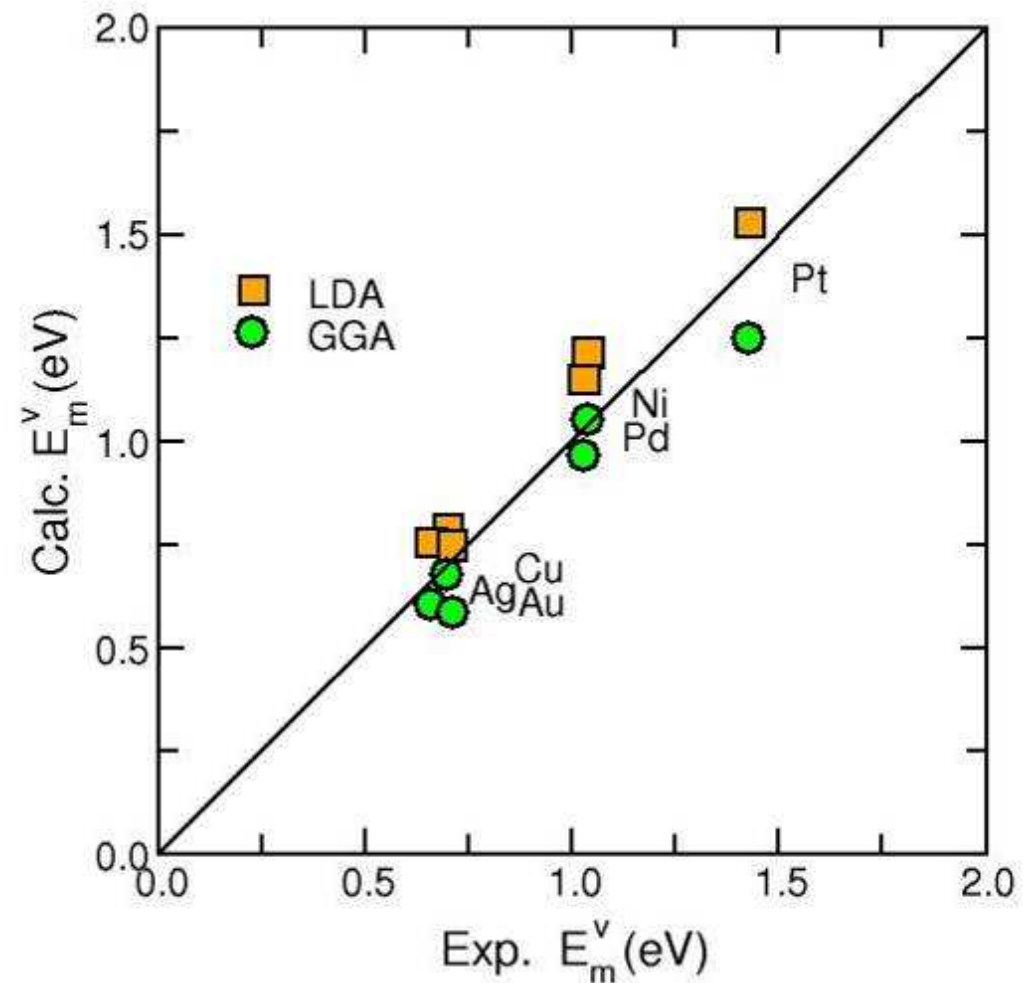


GGA values lower than LDA
by $\sim 10\%$

Vacancy formation energy



Vacancy migration energy



- LDA in better agreement with exp.
- GGA underestimates vac. form. energy

$GGA < \text{Expt.} < LDA$

$$E_{\text{tot}} = E_{\text{bande}} + E_{\text{rep}}$$

Terme de bande :

- $$E_{\text{bande}} = \sum_n^{\text{occupés}} \langle \psi_n | \mathbf{H}(\{r_{ij}\}) | \psi_n \rangle$$
- pour construire la matrice $5N \times 5N$ $H_{i\alpha, j\beta}$ (métaux de transition) ; 2 sortes de paramètres
- pour les termes *non-diagonaux*, les paramètres de Slater-Koster pour les intégrales de saut :
 $V_{dd\sigma}(r), V_{dd\pi}(r), V_{dd\delta}(r)$
de type V_0 *fonction décroissante de r
- pour les termes *diagonaux*, les énergies de site : $e_{s,i}, e_{p,i}$, constantes, ou dépendantes de l'environnement i.e. $\{r_{ij}\}$

Terme répulsif :

- potentiel de paires :
$$E_{\text{rep}} = \sum_{ij} \phi(r_{ij})$$
- ou plus général :
$$E_{\text{rep}} = \sum_i f \left(\sum_j \phi(r_{ij}) \right)$$

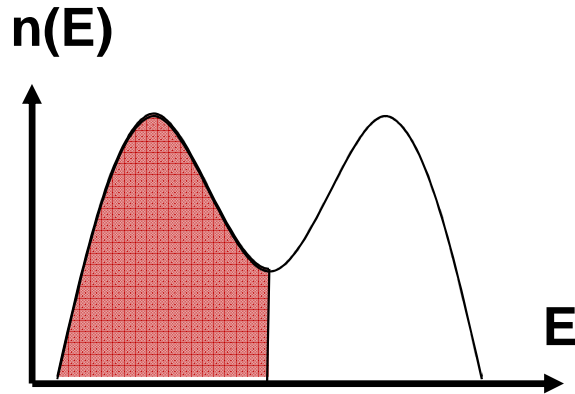
En pratique :

- 4 à 100 paramètres ajustables
- ajustement sur l'expérience ou les calculs ab initio
(structure de bande, courbes énergie/volume)

Remarque :

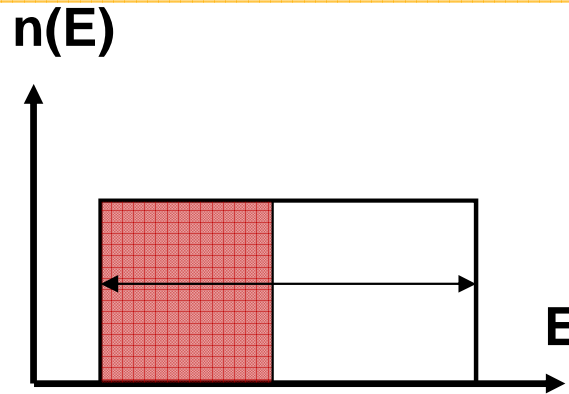
- c'est la forme la plus simple : orthogonale à 2 centres
- il existe de formes plus compliquées : non orthogonales et à plusieurs centres

Liaisons Fortes



- Forces angulaires
- Propriétés dépendent du remplissage de la bande d

Approximation du Second moment



$$E^i = E_{\text{paire}}^i + E_{\text{bande}}^i$$

$$E_{\text{bande}}^i = \sqrt{\rho_i}$$

- Propriétés normalisées indépendantes du remplissage
- Pas de forces angulaires
- 4 paramètres ajustés sur données expérimentales (E_{coh} , C_{ij} , a)

Potentiel EAM

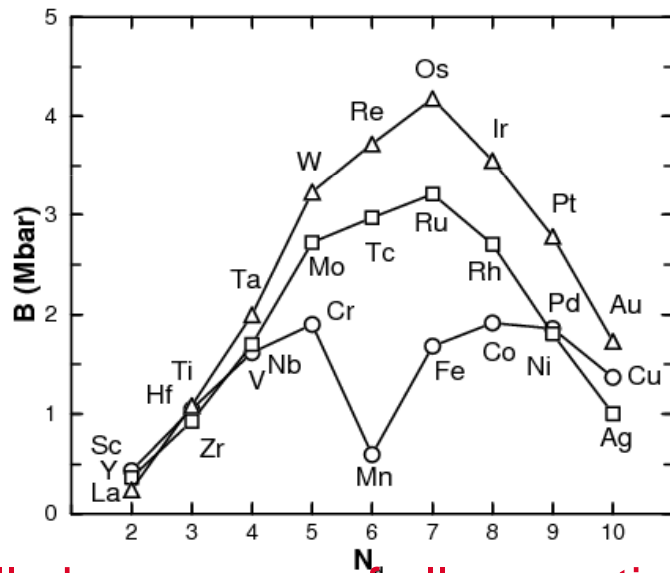
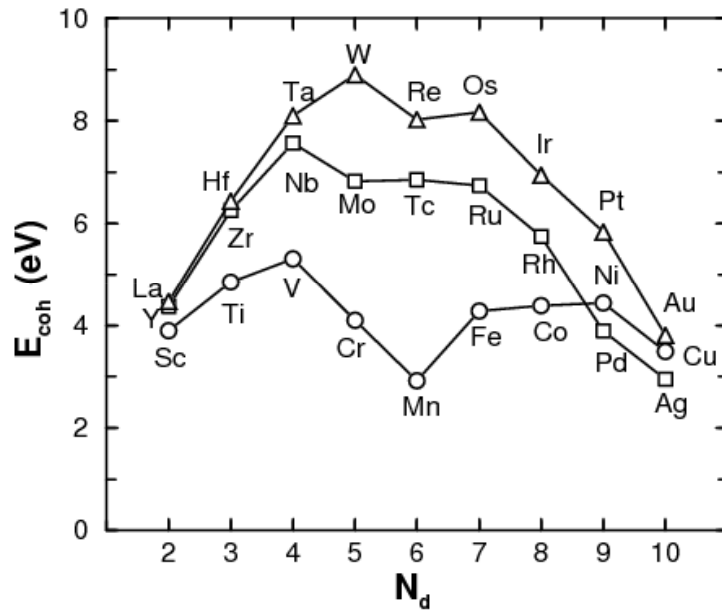
$$E^i = E_{\text{paire}}^i + F(\rho_i)$$

Mendelev et al. (2003):

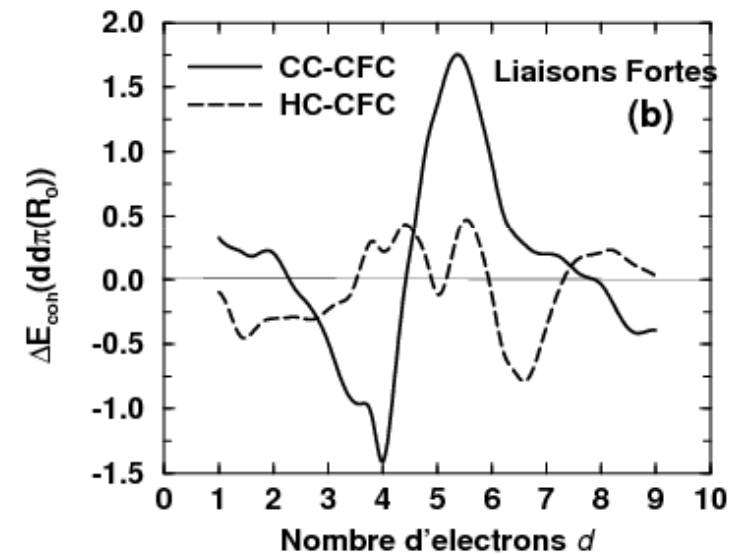
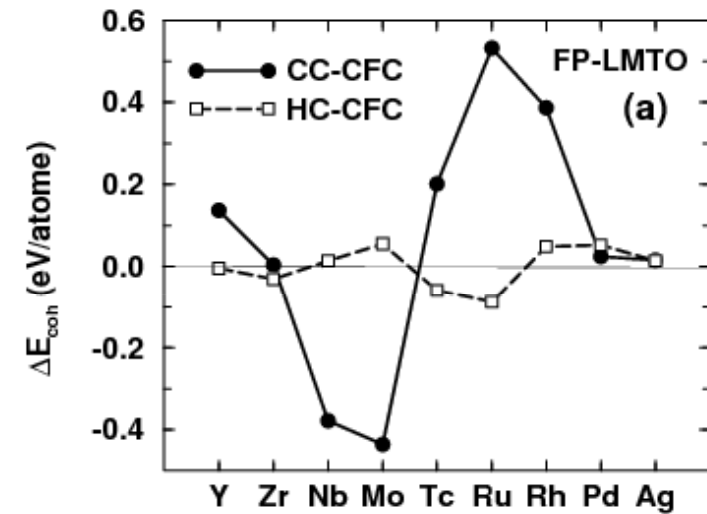
$$F(\rho_i) = \sqrt{\rho_i} + a\rho_i^2$$

- = généralisation empirique du second-moment
- + de flexibilité dans dépendance en distance
- Ajustement des paramètres (10-15) sur résultats DFT (défauts, "liquide")

VARIATIONS ALONG TRANSITION METAL SERIES



Bell-shape curve of all properties:
second-moment or band filling effect

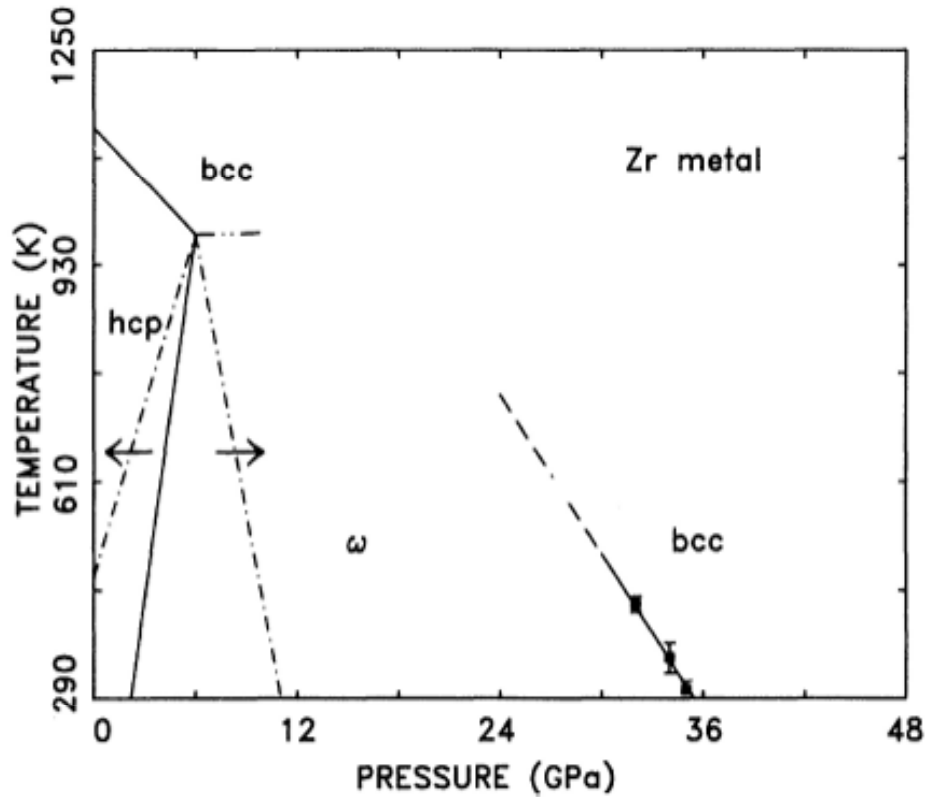


Relative stabilities are sensitive to
details of the density of states: well
reproduced by TB d-band model | PAGE 19

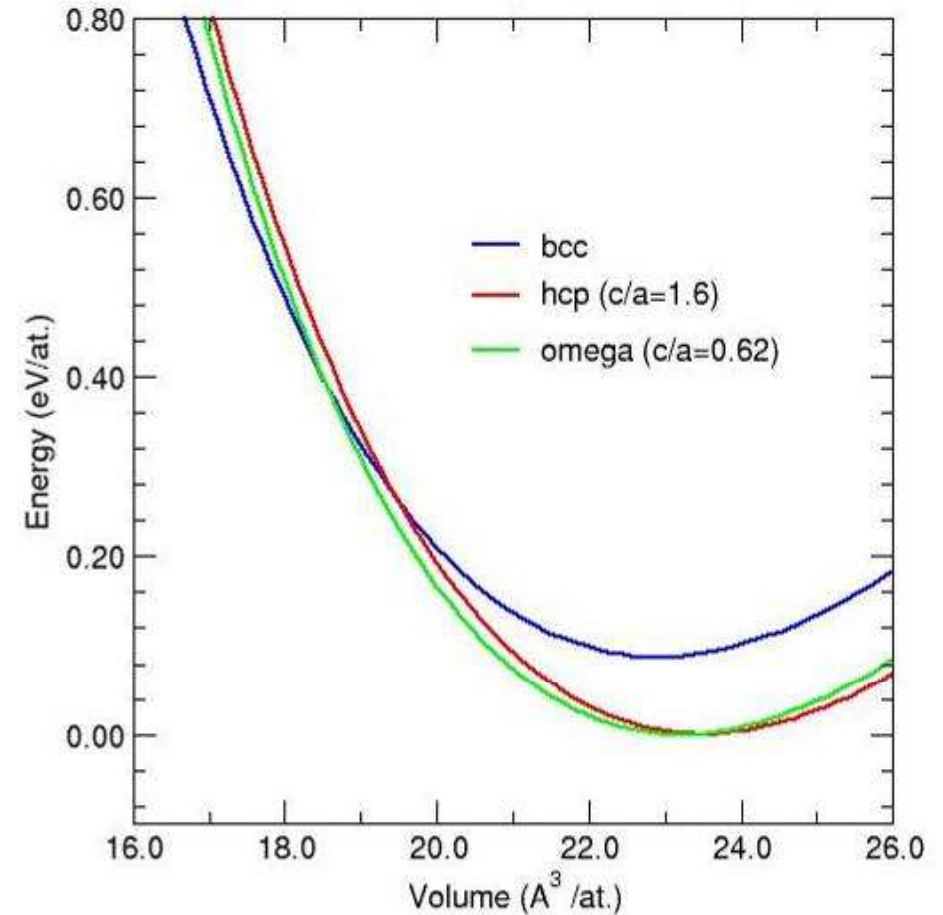
	Temps CPU 250 atomes 1 configuration	Variation avec le nombre d'atomes (N)	Avantages et domaine de validité	Inconvénients et limitations
Potentiel empirique	10^{-3} sec	N	<ul style="list-style-type: none"> • études génériques • Valide près du domaine d'ajustement 	<ul style="list-style-type: none"> • paramètres à ajuster • alliages très difficiles • ΔE approximatif
Liaisons Fortes	10 mn	N (récursion) à N^3 (diagonalisation)	<ul style="list-style-type: none"> • contient toute la physique • Très bon pour les tendances • Propriétés électroniques • Energies par atome 	<ul style="list-style-type: none"> • paramètres à ajuster • alliages difficiles
DFT bases localisées (SIESTA)	6 heures	N^3	<ul style="list-style-type: none"> • Énergies quantitatives • Alliages 	<ul style="list-style-type: none"> • Bases à développer
DFT ondes planes	15 jours	N^3	<ul style="list-style-type: none"> • très robuste 	<ul style="list-style-type: none"> • temps de calcul

EXAMPLES OF DFT CALCULATIONS OF BASIC QUANTITIES

Experimental phase diagram



Relative stability at T=0
DFT-GGA



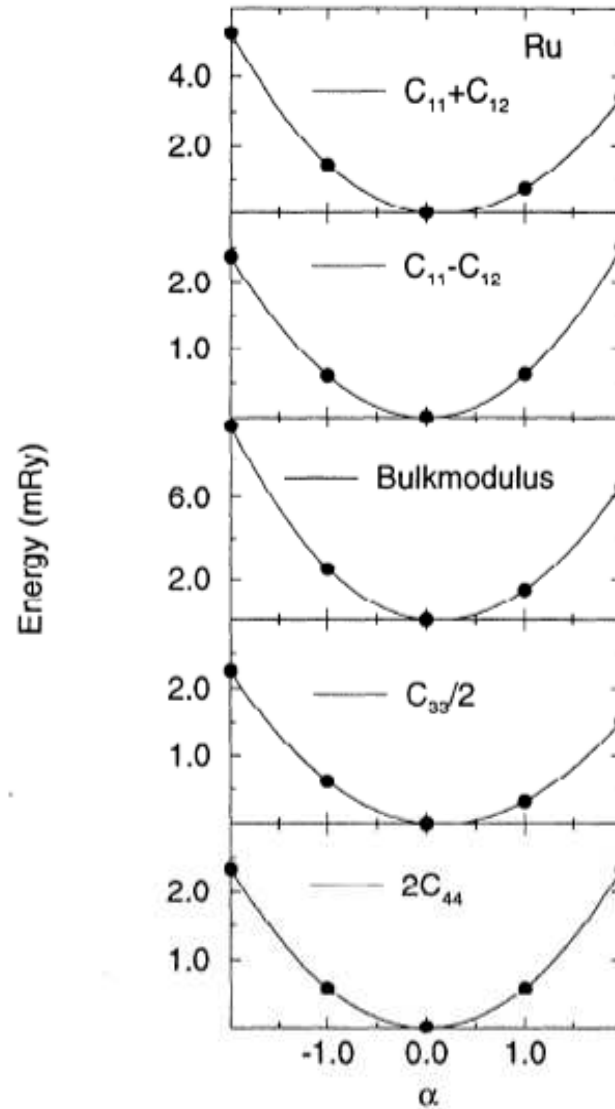
- Predicts successive phases under pressure: hcp- ω -bcc
- $P_{\omega\text{-bcc}} = 30 \text{ GPa}$

$$\begin{pmatrix} 1+\alpha & 0 & 0 \\ 0 & 1+\alpha & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\begin{pmatrix} 1+\alpha & 0 & 0 \\ 0 & 1-\alpha & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1+\alpha \end{pmatrix}$$

$$\begin{pmatrix} 1 & 0 & \alpha \\ 0 & 1 & 0 \\ \alpha & 0 & 1 \end{pmatrix}$$



hcp Ru

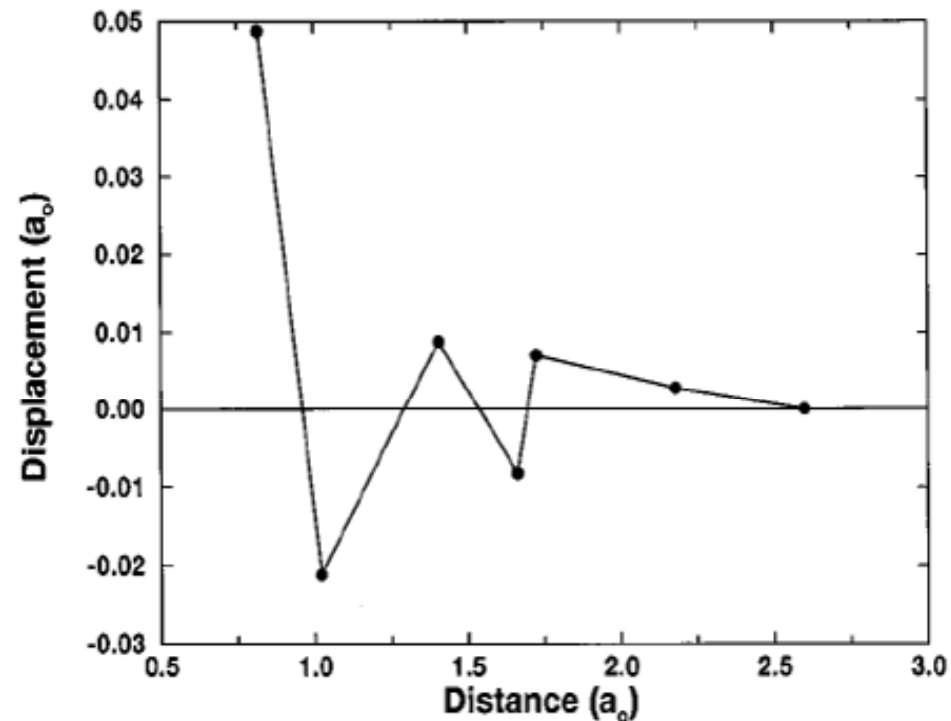
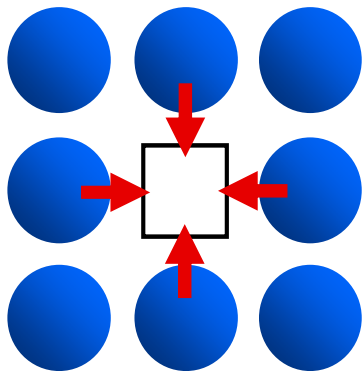
Accord avec
l'expérience des
calculs DFT :
10 à 20%

FIG. 1. Total energy of Ru as function of the lattice distortion parameter α , for the five different strains defined in the text. From the top panel to the bottom panel, the strains correspond to the distortions given in Eq. (5), Eq. (7), Eq. (14), Eq. (9), and Eq. (11), respectively.

L. Fast et al.
(1995)

« Relaxation » or geometry optimisation = minimization of the energy, $E(\{\mathbf{r}_i\})$, wrt atomic positions, $\{\mathbf{r}_i\}$.

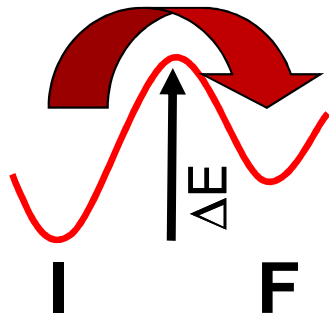
Example : around a vacancy



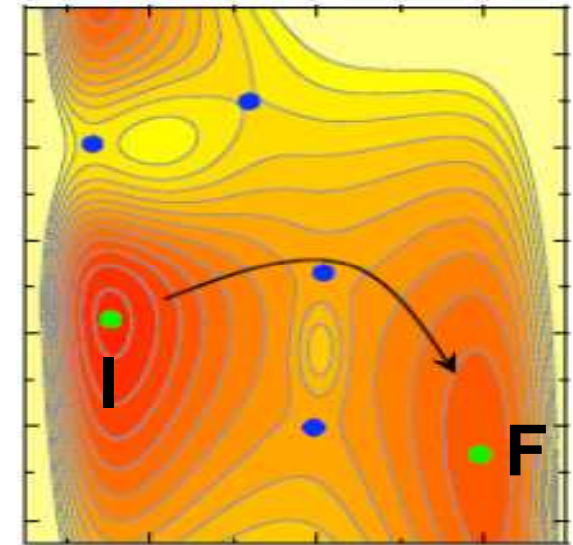
In Ta, relaxation reduces the DFT value of vacancy formation energy from 3.5 eV to 3.0 eV (Expt. = 3.1 eV)

A. Satta et al.
Phys. Rev.B (1999)

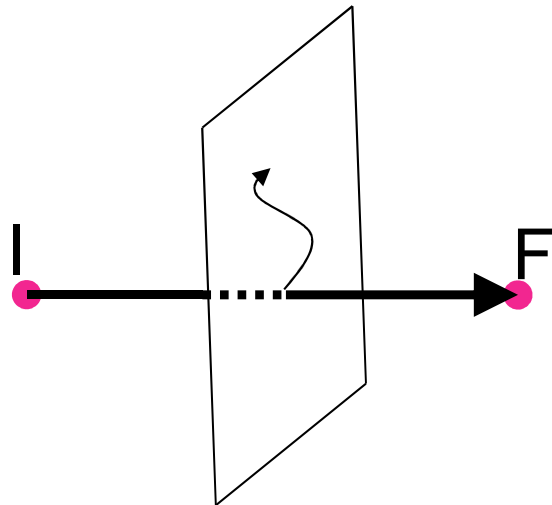
- Key ingredient for thermally activated processes:



$$\Gamma(T) = \Gamma_0 \exp\left(-\frac{\Delta E}{k_B T}\right)$$



- reaction coordinate method:



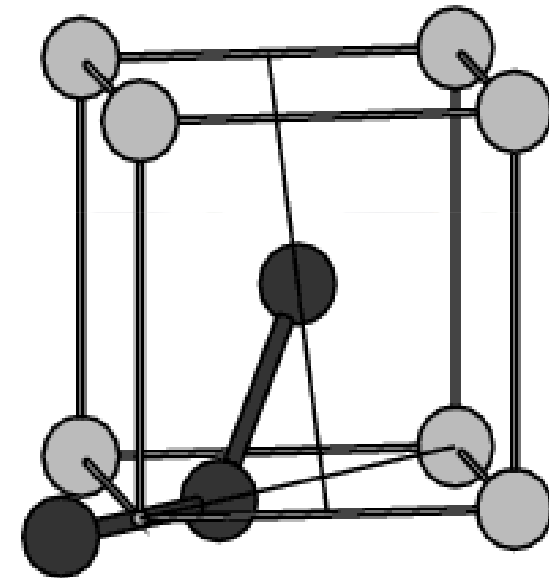
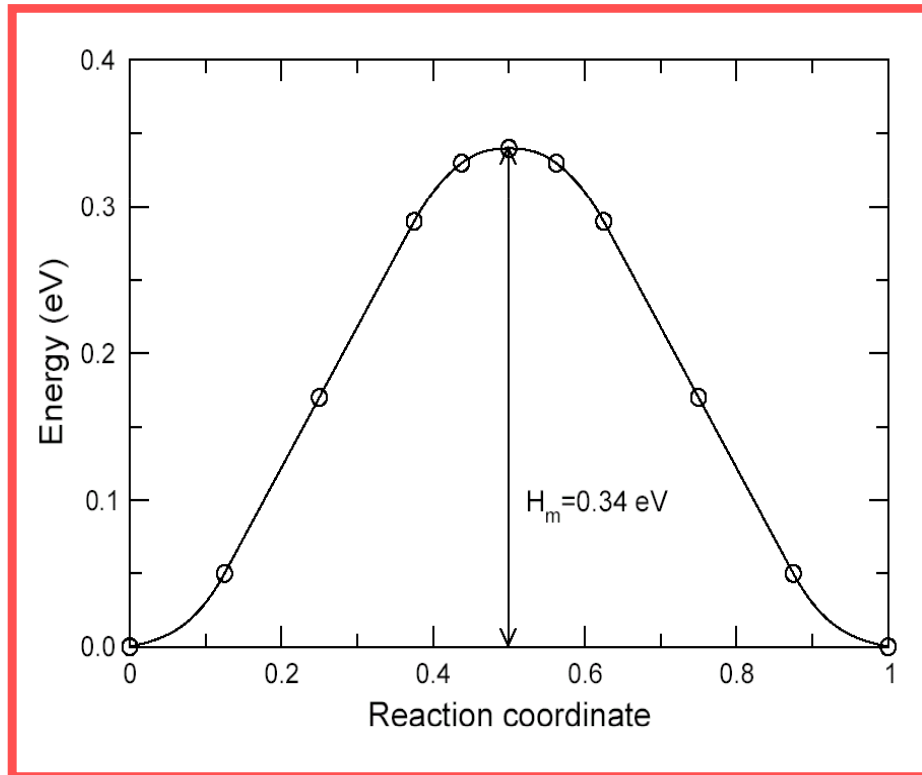
- Minimization with a constrain, e.g.:

$$\xi(R) = \frac{(R - R_{CM})(R_F - R_I)}{|R_F - R_I|^2}$$

relaxation in the hyperplan perpendicular to the direct path

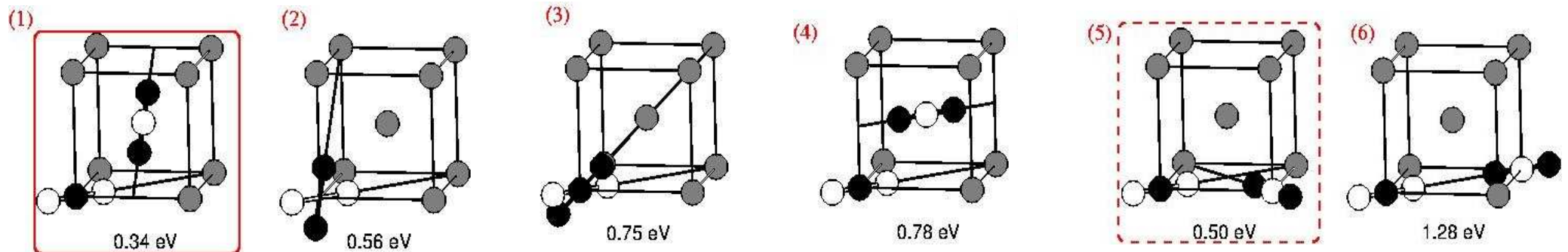
- Other method: Nudged Elastic Band (NEB)

Self-interstitial in Fe (DFT)

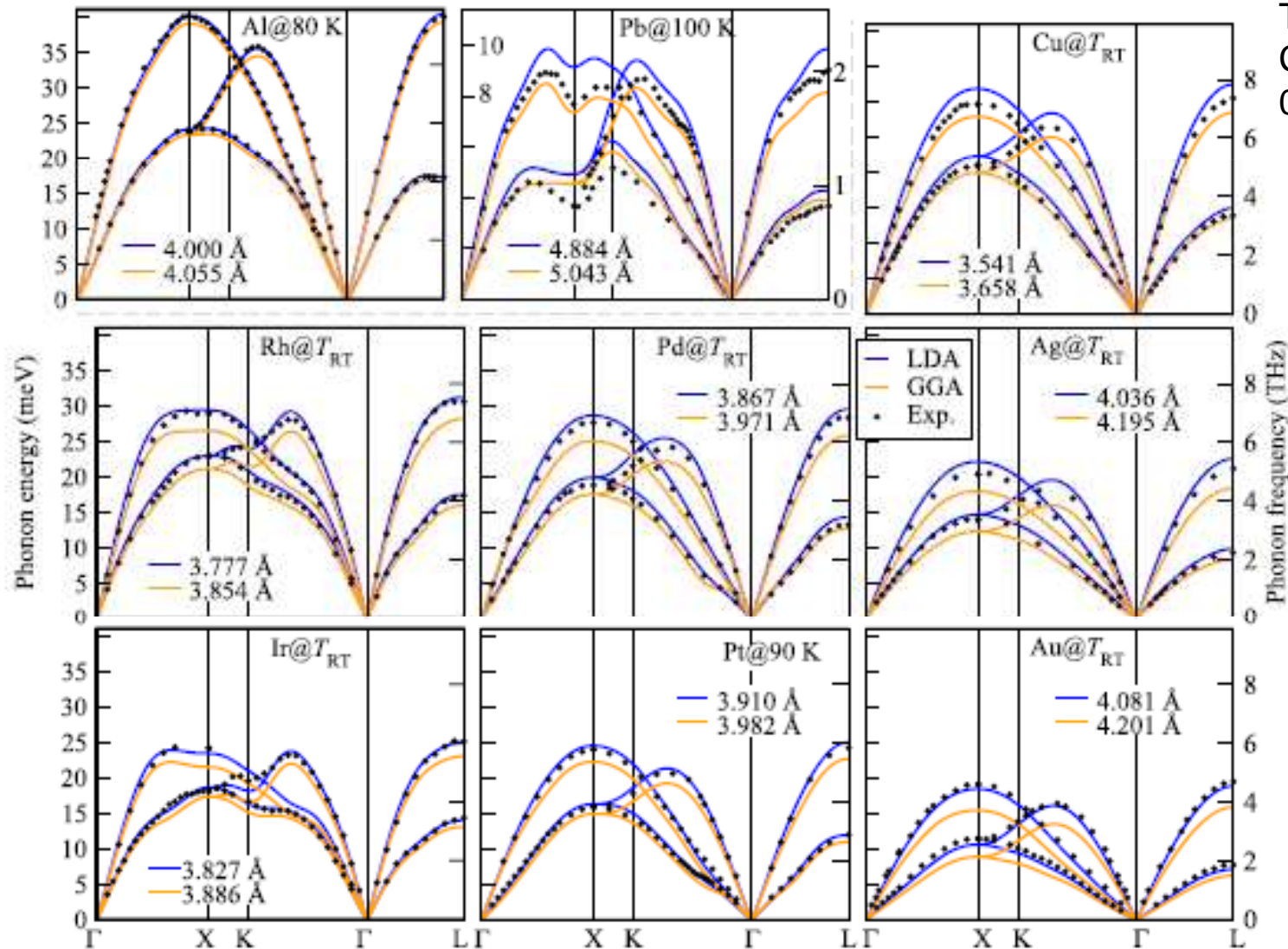


- Excellent agreement with exp. : $H_M = 0.3 \text{ eV}$

C.C. Fu et al. Phys. Rev. Lett (2002)



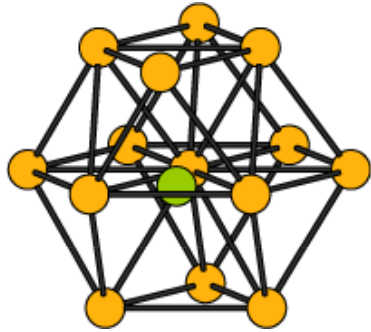
- Bulk phonons:



T. Hickel, et al. J. Phys.:
Condens. Matter 24 (2012)
053202

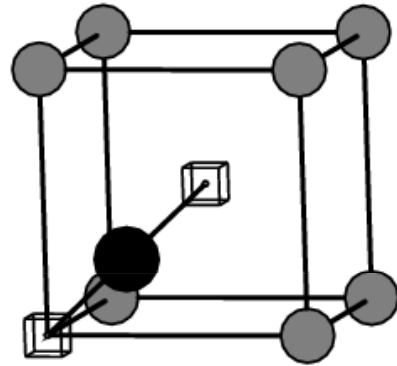
- Can be used to compute defect attempt jump frequencies (Transition state theory)

EXAMPLES OF RECENT DFT STUDIES



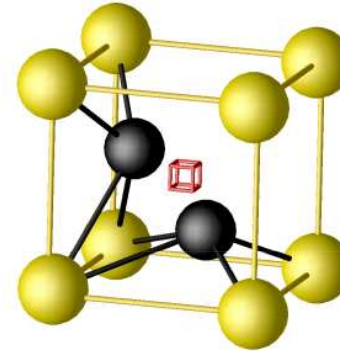
Self-interstitials in hcp metals

Vérité (2007)
Gasca (2010)



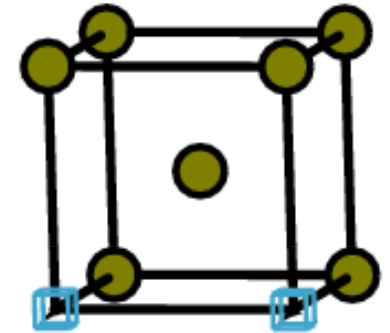
Helium-vacancy complexes in Fe

Fu (2005)



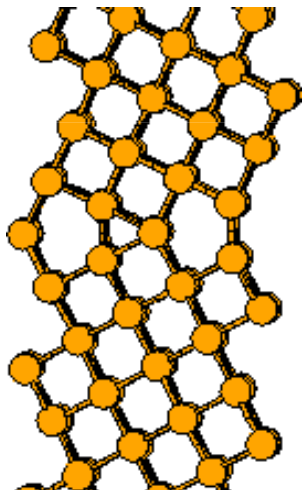
Vacancy-carbon complexes in Fe

Domain (2004)
Fu (2007)



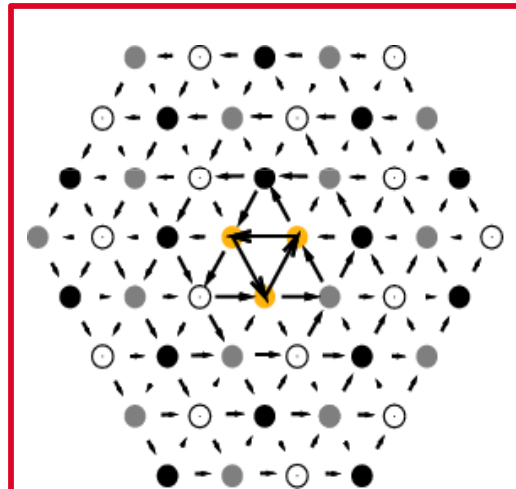
Di-vacancies in bcc metals

Ventelon (2012)



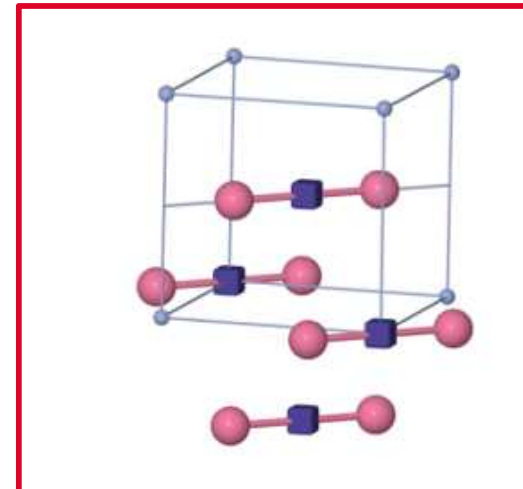
Grain boundary in Fe

Fu (2011)



Dislocation cores in bcc metals

Ventelon (2012)

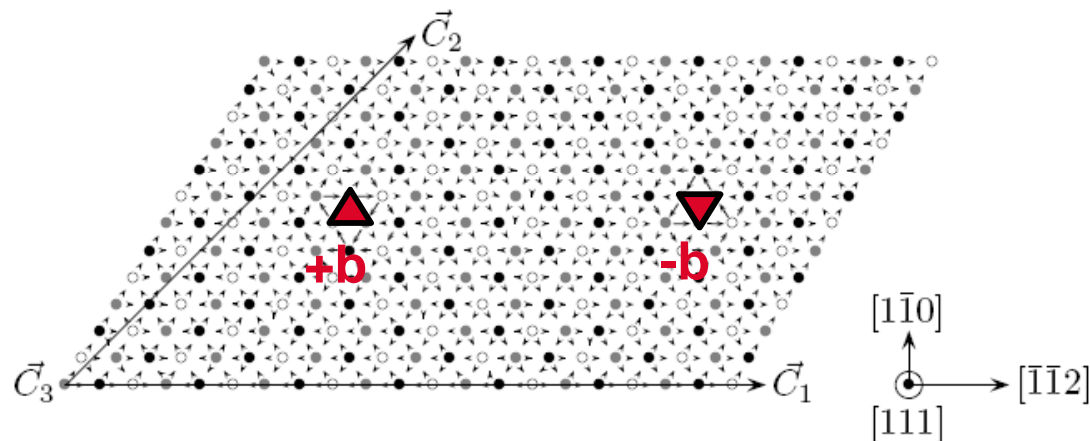


Self-interstitial clusters in Fe

Marinica (2012)

CORE OF $\langle 111 \rangle$ SCREW DISLOCATIONS IN BCC METALS (DFT)

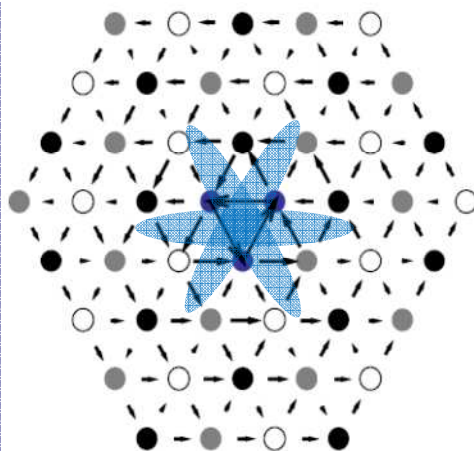
Simulation cell: 273 atom cell



Core structure

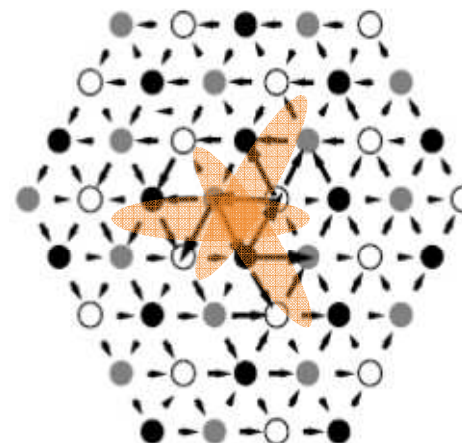
- Vitek representation
- Most empirical potentials predict degenerate cores

Fe, Mo, Ta, W



Non degenerate core

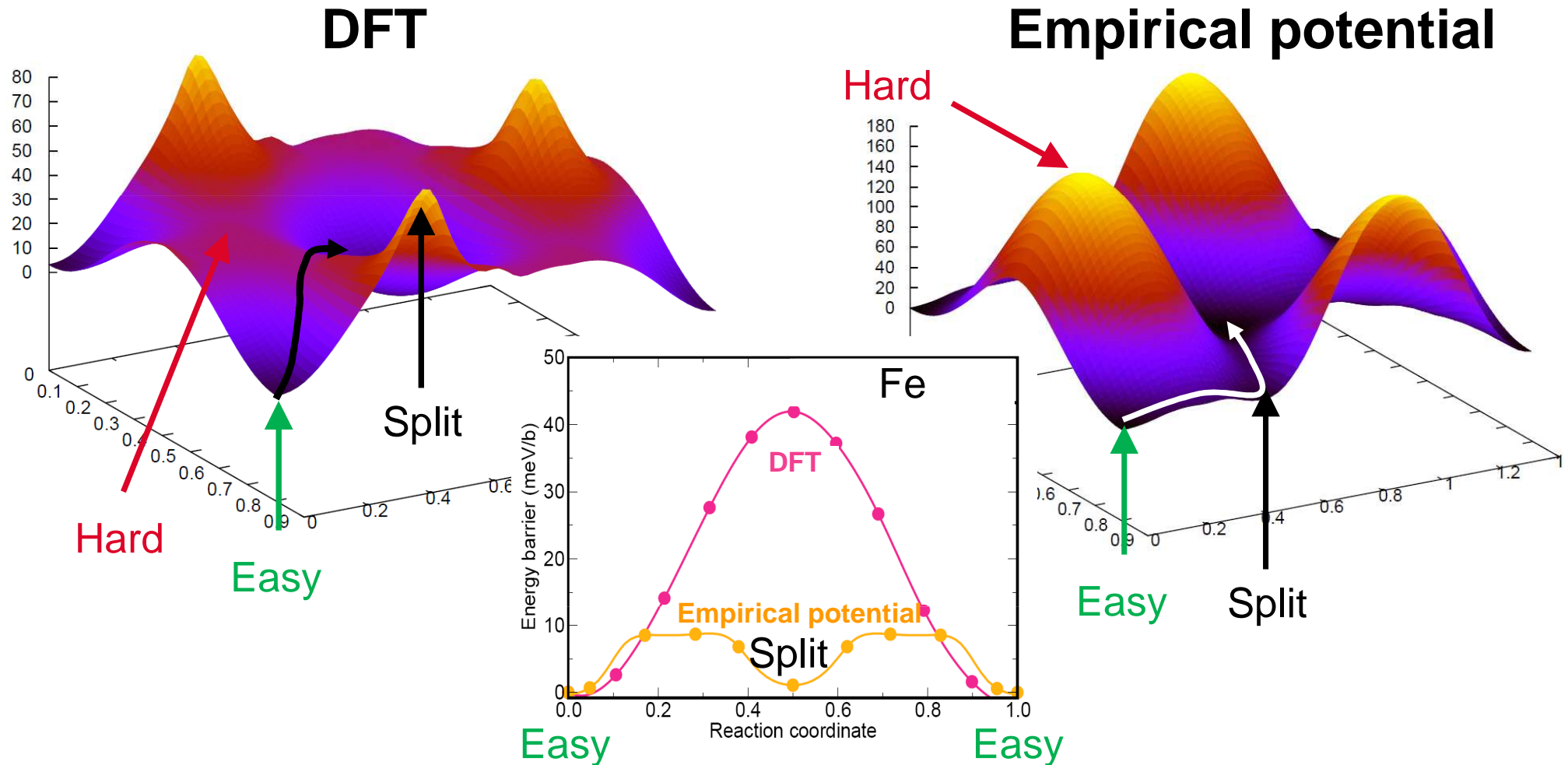
W-25%Re



Degenerate core

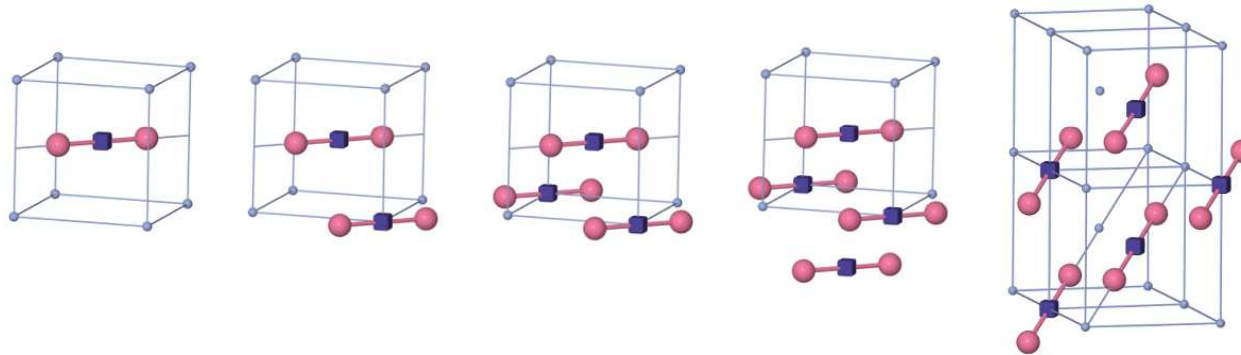
Romaner (2010)
Ventelon (2011)

2D PEIERLS POTENTIAL OF <111> SCREW DISLOCATIONS IN BCC FE



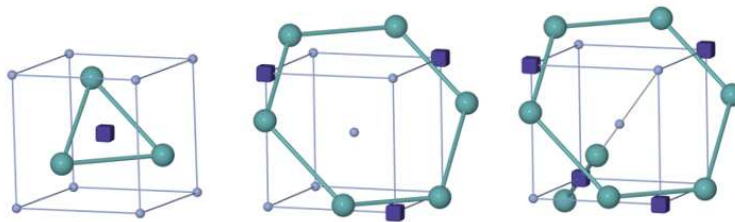
Compatible with absence of asymmetry in the crystal orientation dependence of the critical flow stress observed experimentally in iron at low temperature

SELF INTERSTITIAL CLUSTERS IN BCC FE



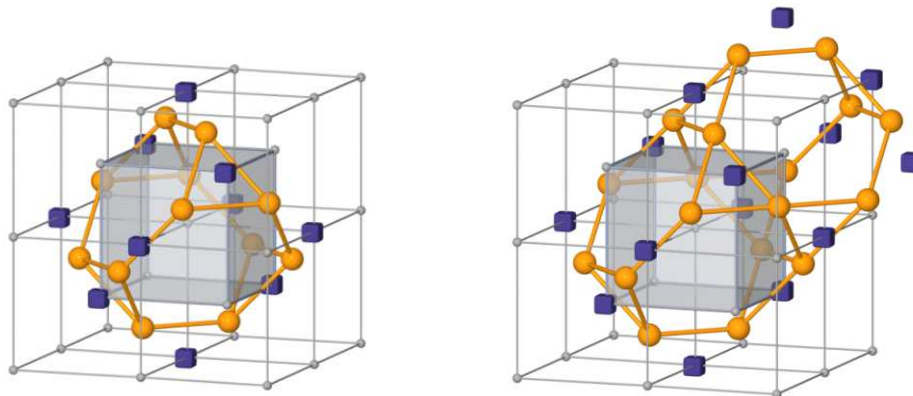
Parallel dumbbells

FW et al.; NIMB 228 (2005) 92



Triangle and rings

D.A. Terentyev et al.
PRL 100 (2008) 145503



3D structure
C15 nanocrystals

M. C. Marinica, FW, J-P Crocombette
Phys. Rev. Lett.108 (2012) 025501

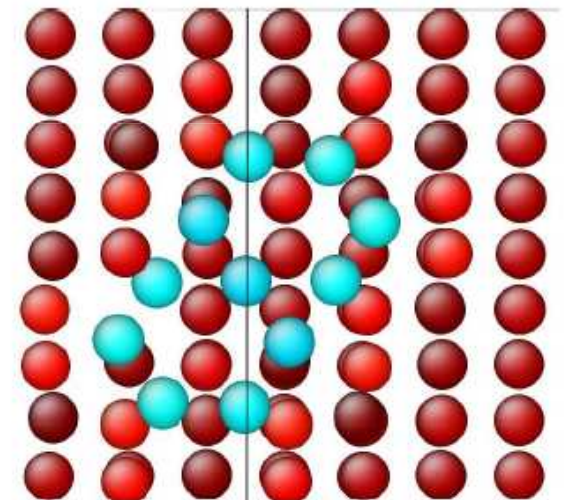
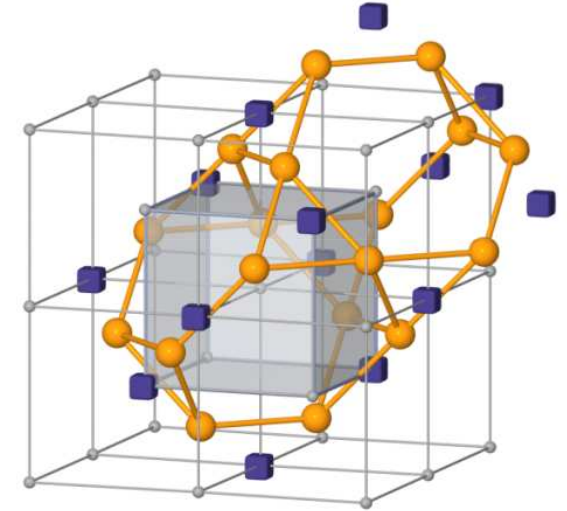
3D crystalline structure predicted for self-interstitial clusters in bcc metals

DFT results:

- Very low energy structures in Fe, and to a lesser extend in Ta (DFT)
- They have large antiferromagnetic moments

Empirical potential (fitted to DFT):

- They are formed in cascades
- They can grow by capturing $\langle 110 \rangle$ dumbbells
- They are immobile



DFT

- exchange-correlation functionals
- parallelization beyond 1000 cores
- finite temperature effects in DFT
- toolkits (phase diagrams, Onsager coefficients, etc.)
- automatic sampling of energy landscapes (ART, dimer)
- dislocation kinks

Other atomistic methods

- Empirical potentials for alloys and tight-binding models (eg Fe-C)
- Slow dynamics codes (CRESCENDO CD code)

Coupling with experiments

- Electronic Structure: Basic Theory and Practical Methods, R. M. Martin (Cambridge University Press, 2004)
- Handbook of Materials Modeling, Ed. S. Yip (Springer, 2005)
- Comprehensive Nuclear Materials, Vol. 1: Basic Aspects of Radiation Effects in Solids/Basic Aspects of Multi-Scale Modeling, Ed. R. J.M. Konings (Elsevier, 2012)

